

Composite Higgs searches

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'Next steps in the Energy Frontier - Hadron Colliders' workshop

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Introduction

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The LHC is doing a wonderful job in testing the TeV-scale physics

After the full LHC program we will gain a lot of information

- fair determination of single **Higgs observables** ($\sim 10\%$)
- test of **naturalness** up to $\mathcal{O}(\text{few}\%)$

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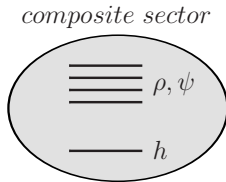
➤ Which is the situation in composite Higgs models?

Introduction: the composite Higgs scenario

In composite Higgs models the Higgs arises as a bound state of a **new strongly-coupled sector**

The composite dynamics gives rise to additional states at the TeV scale:

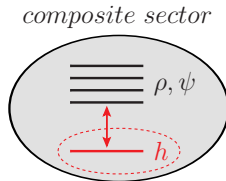
- Fermionic top partners
- Spin-1 resonances
(KK-gluons and EW resonances)



Introduction: the composite Higgs scenario

Need **mass gap** between the Higgs and the other composite states

➤ natural if the Higgs is a **Goldstone**



Key features:

- Modification of the SM Higgs couplings
- Non-linear dynamics ➡ new non-renormalizable interactions

Introduction

Indirect bounds and lack of enough luminosity limit some important searches at the LHC:

- EW precision data strongly favor small deviations in the **linear Higgs couplings** ($< 10\%$)
- EW precision data push the mass of **vector resonances** to the multi-TeV range
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 - tests of the **non-linear Higgs dynamics** usually rely on rare processes with small cross sections
- What kind of new experiments do we need to extend our reach on composite models?

Introduction

In this talk I will discuss how much a future 100 TeV hadronic collider could improve our reach on the parameter space of composite Higgs scenarios:

- ▶ testing the Higgs compositeness
- ▶ searching for top partners
- ▶ searching for vector resonances

Testing Higgs compositeness

Higgs compositeness

The **Goldstone boson nature** of the Higgs implies a well defined pattern of deformations of the Higgs couplings

$$\mathcal{L} = m_W^2 W_\mu^+ W^{-\mu} \left(1 + 2 \kappa_v \frac{h}{v} \right) - \sum_\psi m_\psi \bar{\psi} \psi \left(1 + \kappa_f \frac{h}{v} \right) + h.c.$$

❖ size of the corrections controlled by the compositeness scale f

- The couplings to the gauge fields only depend on the Goldstone structure (eg. $SO(5) \rightarrow SO(4)$)

$$\text{MCHM}_4, \text{MCHM}_5 \quad \kappa_v = \sqrt{1 - \xi} \quad \xi = v^2/f^2$$

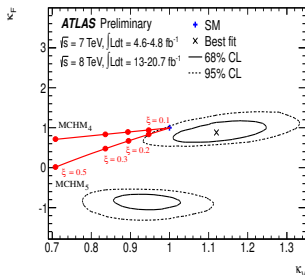
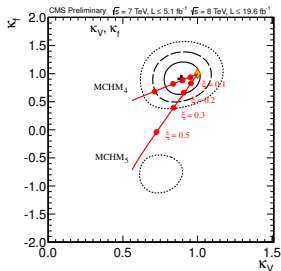
- The couplings to the fermions have more model dependence

$$\begin{array}{ll} \text{MCHM}_4 & k_f = \sqrt{1 - \xi} \\ \text{MCHM}_5 & k_f = \frac{1 - 2\xi}{\sqrt{1 - \xi}} \end{array} \quad \xi = v^2/f^2$$

Single Higgs couplings

Measuring κ_ν gives a **model-independent** bound on ξ

The current data give a bound $\xi \lesssim 0.2$



- The ultimate LHC reach is $\xi \sim 0.1$
- **Limited improvement** possible at an hadronic machine!

Higgs couplings at a 100 TeV machine

Other measurements can however benefit from high energy and high luminosity:

- Higgs coupling to the **top quark**
- **non-linear** Higgs interactions

Possible relevant channels:

- WW scattering and $WW \rightarrow hh$
- Double Higgs production in gluon fusion

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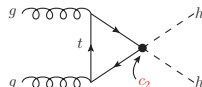
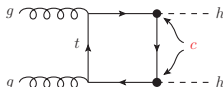
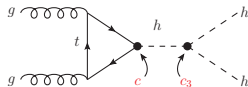
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Double Higgs production in gluon fusion

The relevant Higgs couplings can be parametrized as

$$m_{top} \bar{t} t \left(c \frac{h}{v} + \frac{c_2}{2} \frac{h^2}{v^2} \right), \quad c_3 g_{hhh}^{SM} h^3$$



➤ can be used to extract the Higgs trilinear coupling c_3

[Baur, Plehn, Rainwater; Grober, Muhlleitner;

Contino, Ghezzi, Moretti, G. P., Piccinini, Wulzer;

Dolan, Englert, Spannowsky; Goertz, Papaefstathiou, Yang, Zurita; ...]

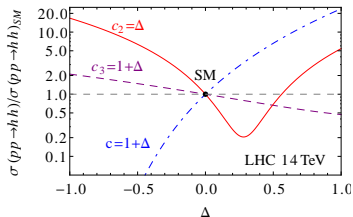
➤ is sensitive to **non-renormalizable Higgs interactions**
($\bar{t} t h h$ is a distinctive sign of a composite Higgs)

[Grober, Muhlleitner; Contino, Ghezzi, Moretti, G. P., Piccinini, Wulzer]

Double Higgs production in gluon fusion

The cross section can be significantly **modified** even for small deviations of the Higgs couplings

- strong dependence on c and c_2
- milder dependence on c_3



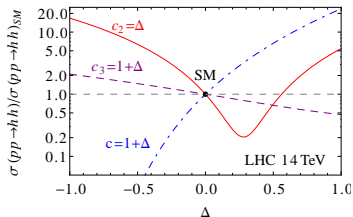
At the LHC the measurement is severely limited by statistics:

$$\sigma(gg \rightarrow hh \rightarrow \gamma\gamma bb)_{SM} \simeq 5 \text{ ab @ 14 TeV}$$

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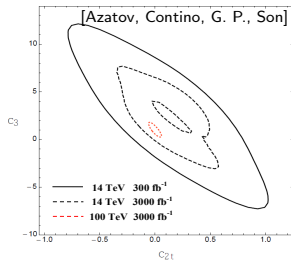


At the LHC the measurement is severely limited by statistics:

$$\sigma(gg \rightarrow hh \rightarrow \gamma\gamma bb)_{SM} \simeq 5 \text{ ab @ 14 TeV} \quad \Rightarrow \quad 200 \text{ ab @ 100 TeV}$$

- **huge improvement** at 100 TeV!

Can test Higgs trilinear
and Higgs compositeness for $\xi \sim 0.02$



Searches for top partners

Top partners and naturalness

Top partners control the generation of the Higgs potential and the stability of the Higgs mass

$$\delta m_h^2|_{1-loop} \sim \text{diagram 1} + \text{diagram 2} \sim -\frac{y_{top}^2}{8\pi^2} M_X^2 \lesssim \text{TeV}$$

The diagram shows two Feynman diagrams for the one-loop correction to the Higgs mass. The first diagram is a loop of top quarks, with external Higgs lines labeled 'h'. The second diagram is a loop of a new particle labeled 'NP', also with external Higgs lines labeled 'h'. The NP loop is shaded gray.

A **bound** on the partners mass translates into an unavoidable amount of **fine-tuning**

$$\Delta \gtrsim \left(\frac{M_X}{400 \text{ GeV}} \right)^2$$

Top partners and naturalness

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The diagram shows two Feynman diagrams for the one-loop correction to the Higgs mass. The first diagram is a loop of top quarks, with external Higgs lines labeled 'h' and the loop labeled 'top'. The second diagram is a loop of a new particle labeled 'NP', with external Higgs lines labeled 'h'.

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Natural Composite Higgs:
light top partners

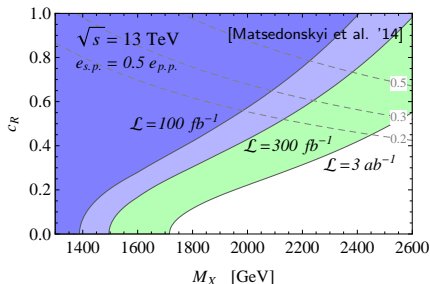


Natural SUSY:
light stops

Top partners at the LHC

Top partners can be easily produced through:

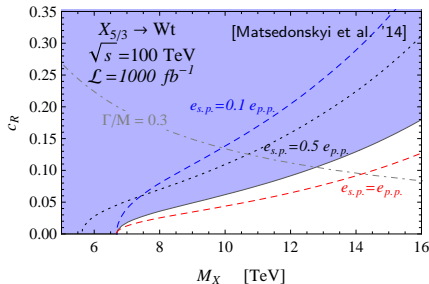
- QCD interactions (pair prod. – universal)
- couplings with heavy quarks (single prod. – model dependent)



- LHC can probe partner masses up to $M_X \sim 2 \text{ TeV}$
- testing **naturalness** up to $\mathcal{O}(4\%)$

Top partners at 100 TeV

At 100 TeV the production cross section allows to explore a huge range of masses



- model-independent reach $M_X \sim 7 \text{ TeV}$
- by using single production $M_X \sim 12 \text{ TeV}$
- testing **naturalness** up to $\mathcal{O}(0.1\%)$

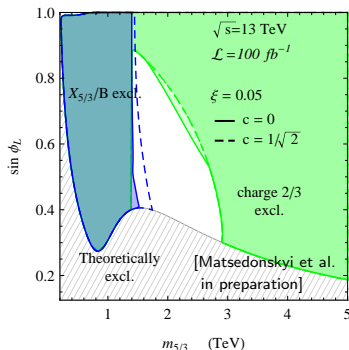
Implications on explicit models

In a large class of models (eg. MCHM_{4,5,10}) the mass of the lightest partner is tightly connected to the compositeness scale

[Matsedonskyi, G. P., Wulzer; Marzocca, Serone, Shu; Pomarol, Riva]

$$\frac{m_H}{m_{top}} \gtrsim \frac{\sqrt{3}}{\pi} \frac{M_X}{f} \quad \Rightarrow \quad \xi \lesssim \left(\frac{500 \text{ GeV}}{M_X} \right)^2$$

- LHC can test $\xi \sim 0.05$
- at 100 TeV can probe $\xi \sim 0.002$



Searches for vector resonances

Vector resonances

Vector resonances are unavoidably present in composite Higgs models

The composite dynamics is charged under an extension of the SM gauge groups

- QCD vector resonances (“KK gluons”)
- EW vector resonances

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EW vector resonances

EW vector resonances are mixed with the SM gauge bosons and generate corrections to the EW observables

$$\Delta\hat{S} \simeq \frac{m_W^2}{m_\rho^2}$$

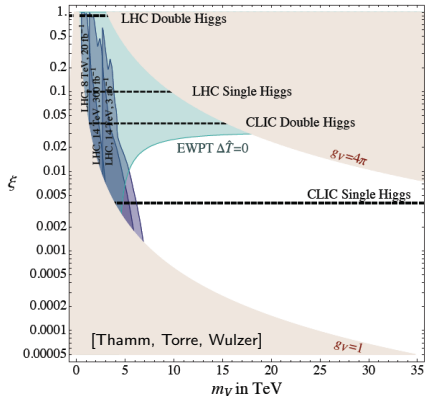
The EW precision data put strong (indirect) bounds

$$m_\rho \gtrsim 2 \text{ TeV}$$

- Much stronger bound if no tuning in \hat{T} : $m_\rho \gtrsim 5 \text{ TeV}$

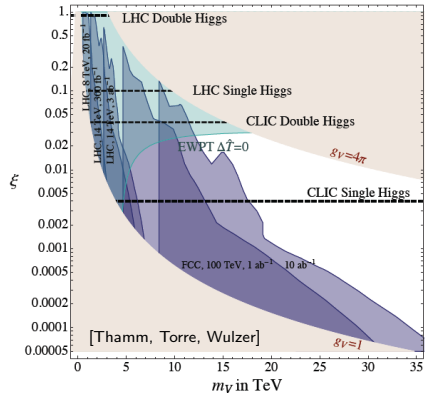
EW vector resonances

- most of the parameter space accessible at the LHC is already **disfavored** by the EW data



EW vector resonances

- ▶ most of the parameter space accessible at the LHC is already **disfavored** by the EW data
- ▶ an 100 TeV collider can easily probe regions with **tiny Higgs compositeness** ($\xi \lesssim 0.002$)



Conclusions

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Thanks to the large increase in energy and cross sections a future 100 TeV hadronic collider can hugely improve the searches for composite Higgs scenarios

- Significantly **extend** LHC searches
 - reach top partners with $M_X \sim 10$ TeV
 - test naturalness up to $\mathcal{O}(0.1\%)$
- **Probe** observables and dynamics not testable at the LHC
 - non-linear Higgs couplings
 - EW vector resonances